

Time-resolved resonant soft x-ray diffraction with free-electron lasers: Femtosecond dynamics across the Verwey transition in magnetite

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Resonant soft x-ray diffraction (RSXD) with femtosecond (fs) time resolution is a powerful tool for disentangling the interplay between different degrees of freedom in strongly correlated electron materials. It allows addressing the coupling of particular degrees of freedom upon an external selective perturbation, e.g., by an optical or infrared laser pulse. Here, we report a time-resolved RSXD experiment from the prototypical correlated electron material magnetite using soft x-ray pulses from the free-electron laser FLASH in Hamburg. We observe ultrafast melting of the charge-orbital order leading to the formation of a transient phase, which has not been observed in equilibrium. © 2011 American Institute of Physics. [doi:10.1063/1.3584855]

Resonant x-ray diffraction, i.e., an x-ray diffraction experiment with the photon energy tuned into resonance with a dipole transition, combines the high spectroscopic sensitivity of x-ray absorption spectroscopy with momentum resolution. Thus it provides not only information on structural aspects of a material but also on its electronic properties. Particularly interesting in this regard are the strong dipole allowed excitations into transition-metal 3d, oxygen 2p, and lanthanide 4f-states, which all occur in the soft x-ray range between 400 and 1600 eV. Within the momentum space that can be probed with resonant soft x-ray diffraction are antiferromagnetic reflections from natural or artificial magnetic structures and superstructure reflections caused by periodic modulations of the electronic state as they occur in a large class of correlated electron systems.¹

A prototype material from the latter is magnetite (Fe₃O₄). Already in the late 1930s Verwey discovered that upon lowering of the temperature below T_V=123 K (the Verwey temperature), magnetite undergoes a first-order phase transition leading to a conductivity decrease by two orders of magnitude.² This Verwey transition involves a change from a cubic inverse spinel high-temperature lattice structure to a complex monoclinic low-temperature phase.³ Verwey proposed as the mechanism behind the drop in electrical conductivity a freezing of the charge fluctuations between octahedrally coordinated (B-site) Fe²⁺ and Fe³⁺ ions

into a charge ordered structure. Calculations further predicted orbital order, i.e., a spatial modulation of the orbital occupation, for the low-temperature phase.⁴⁻⁸ Both charge and orbital order have been observed experimentally,^{3,6,9-11} even though no full consensus exists over all aspects of the order.

One of the superstructure peaks characteristic of the low-temperature phase is (0,0,1/2) (notation refers to the cubic room temperature unit cell with a=8.39 Å). The inset to Fig. 1 shows a scan through this peak along the [001] (L) direction in reciprocal space. In the main panel of Fig. 1 the photon energy dependence of the (0,0,1/2)-peak intensity (red) at the oxygen 1s→2p (K) resonance is compared to the x-ray absorption spectrum (black) from the same sample. Since x-ray absorption probes the whole unoccupied p-derived conduction-band density of states, the energy position of the (0,0,1/2) resonance clearly shows that this peak reflects a spatial modulation of states at the bottom of the conduction band. Local-density approximation (LDA+U) calculations reveal that these states are sensitive to charge/orbital order via O 2p - Fe 3d hybridization.^{7,11} It is these states that also form the energy gap^{12,13} leading to the insulating nature of magnetite below T_V. Hence, the (0,0,1/2) diffraction peak provides us with direct information about the electronic order as well as the energy position of the involved electronic levels. By probing this signal time resolved, we can study the melting of charge/orbital order as well as a change in the band structure as a response to an electronic excitation induced with an infrared laser pulse.

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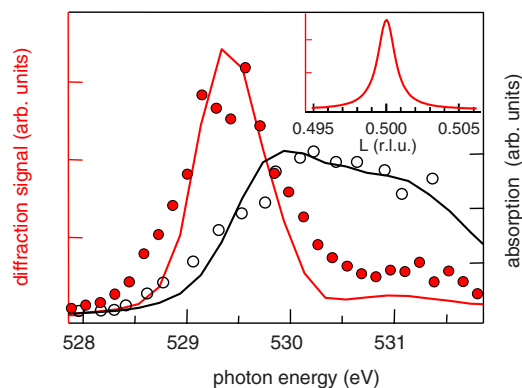


FIG. 1. (Color online) Oxygen K -edge x-ray absorption spectra and (0,0,1/2) RSXD of magnetite ($T=90$ K). The open (filled) symbols represent the fluorescence yield (RSXD signal) recorded at FLASH, the corresponding lines represent the total electron yield (RSXD signal) recorded at BESSY. The inset shows a scan along [001] (L) through the (0,0,1/2) RSXD peak. The FWHM value of 0.0015 r.l.u. (reciprocal lattice units) of the (0,0,1/2) diffraction peak is determined by the x-ray penetration depth.

The time-resolved resonant soft x-ray diffraction (RSXD) measurements were carried out at the PG2 beamline^{14,15} of the free-electron laser facility FLASH at DESY (Hamburg), using monochromatized x-rays of the third undulator harmonic. The monochromator was set to ~ 1 eV photon energy resolution (200 l/mm grating, $c_{\text{gr}}\text{-value}=1.18$) matching the spectral width of the resonant diffraction peak (Fig. 1). Photon energy scans across the Oxygen $1s \rightarrow 2p$ resonance were possible within the FLASH $\sim 1\%$ photon energy bandwidth (i.e., ~ 5 eV).¹⁶ FLASH was operated at a pulse train repetition rate of 5 Hz with each train (macro bunch) containing 30 micro bunches with an internal repetition rate of 500 kHz. The FLASH pulses were typically 40 fs long.¹⁷ When compared to the 500 MHz repetition rate operation of synchrotrons this resulted in similar average photon flux even in the third harmonic of the free-electron laser. In Fig. 1 we compare the spectra recorded at FLASH (symbols) with those recorded at a third generation synchrotron-radiation facility BESSY II in Berlin (lines). Despite the intrinsic intensity fluctuations of a SASE-FEL,¹⁶ spectra with comparable signal-to-noise ratio are obtained at both sources.

A femtosecond (fs) pump-laser¹⁸ ($h\nu=1.5$ eV, 130 fs pulse duration) was synchronized to every other x-ray pulse (250 kHz) allowing for proper normalization by measuring the sample alternating between the ground and the excited state. Zero time delay between infrared pump and x-ray probe pulses was narrowed down to a ~ 20 ps window using an *in situ* high-bandwidth copper photocathode. The pump laser focus on the sample had a ~ 250 μm full width at half maximum (FWHM) diameter providing pump-fluences up to 10 mJ/cm^2 . High-quality synthetic magnetite single crystals were cleaved *in situ* in UHV. Measurements were performed at 90 K sample temperature. The diffracted x-rays were detected by a ~ 3 mm diameter (5° acceptance angle) fast avalanche photodiode (Silicon APD, Laser Components) integrating over the whole diffraction peak. The APD was screened from the pump-laser photons by an optically tight 500 nm thick Al foil. The amplified fast diode signal was directly recorded with a 2 GHz sampling rate resolving each individual micro bunch contribution. The varying temporal jitter between FLASH and fs pump-laser pulses was separately measured for each macro bunch individually using

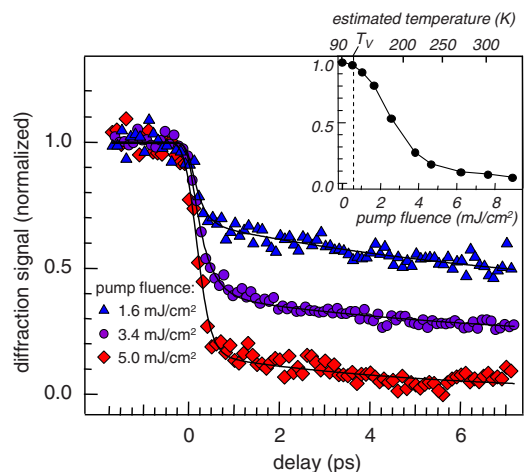


FIG. 2. (Color online) Time-dependent oxygen K -edge (0,0,1/2) RSXD signal vs. laser pump and x-ray probe time-delay for three different fs-laser pulse fluences (symbols). The lines represent double exponential fits. The inset shows the fluence dependence of the (0,0,1/2) RSXD signal measured at 200 ps time delay measured at BESSY II.

electro-optical sampling (EOS).¹⁹ Reference 20 demonstrates that by using EOS for jitter correction a total temporal resolution of 210 fs can be achieved.

In order to study the response of the system to an infrared fs-laser pulse at time t_0 , we monitored the (0,0,1/2) diffraction peak intensity at a second time $t_0 + \Delta t$ using a soft x-ray pulse from FLASH. Figure 2 shows the (0,0,1/2) RSXD peak intensity at 529.4 eV photon energy (peak maximum) as a function of time delay, Δt , for three pump laser fluences 1.6, 3.4, and 5.0 mJ/cm^2 (sample surface). The peak intensity drops by more than 90% within the first ps for the highest pump fluence, while the two lower pump fluences lead to a reduction to $\sim 40\%$ and $\sim 70\%$, respectively, of the initial value. To determine the time constants from the delay scans in Fig. 2, we fitted the data with a double exponential decay convoluted by a Gaussian taking the FLASH temporal resolution into account, which is mainly determined by the arrival time jitter between the two pulses.²⁰ This procedure yields that for all laser fluences the initial fast decay is essentially resolution limited which implies a time constant < 270 fs. In all cases, a subsequent slower decay with smaller amplitudes of $\sim 20\%$ and $\sim 10\%$ follows, the time-scale of which we estimate to several ten ps indicating its relationship with a lattice rearrangement. The inset to Fig. 2 shows the intensity that remains 200 ps after laser excitation as a function of pump fluence measured at BESSY II.

To observe residual (0,0,1/2)-peak intensity 200 ps after the pump pulse is remarkable because the thermalization of the energy deposited by the pump pulse should be completed within a few ps. The sample volume-averaged temperature in the probed sample volume²¹ resulting from this thermalization is shown as the scale on top of the inset of Fig. 2.^{22,23} Except for the weakest, all pump pulses heat the probed sample volume well above $T_V=123$ K. The residual peak intensity observed picoseconds after the pump pulse hence comes from a non-equilibrium transient phase.

In order to characterize the transient phase we make use of the spectroscopic sensitivity of our method. Figure 3 shows a false-color map of the (0,0,1/2) peak intensity as a function of delay time and photon energy for a pump fluence of 4 mJ/cm^2 . The data were recorded by scanning the monochromator and the delay stage for the pump pulses si-

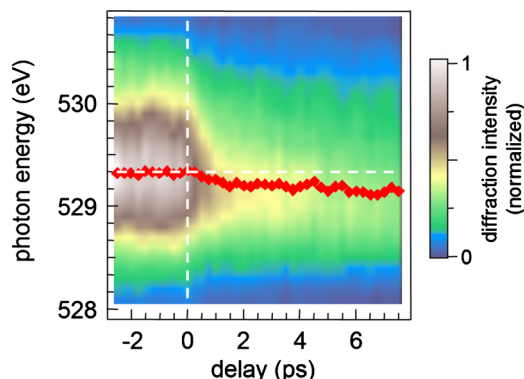


FIG. 3. (Color online) Oxygen *K*-edge (0,0,1/2) RSKD intensity vs. photon energy and time delay. The diamonds mark the resonance peak maxima on the energy scale as determined by individual fits of Gaussian peak profiles to the raw data. The fs laser pump fluence was set to 4 mJ/cm².

multaneously back and forth with sizable different speeds such that the points in the map were recorded in an uncorrelated sequence. This approach averages efficiently over possible long-term drifts in the experimental setup. The temporal resolution for this protracted measurement was 750 fs. We find that the transition from the low-temperature to the transient phase is accompanied by a 0.18 ± 0.05 eV energy shift in the resonance peak to lower photon energies.

The 180 meV size of the energy shift is remarkably large compared to the ~ 40 meV shift observed in x-ray absorption spectroscopy experiments at the Verwey transition.²⁴ In fact the measured shift in the resonance amounts to about the full size of the band gap of 230 meV in the low-temperature phase.⁷ The huge shift in the resonance energy that we observe strongly supports our arguments that a so far unknown transient phase is formed in the first several picoseconds after the pump pulse and we conclude that this transient phase is characterized by the existence of partial charge/orbital order in the presence of an almost closed band gap.

In conclusion we demonstrated how a fs-infrared-laser-pump x-ray-free-electron-laser-probe resonant soft x-ray diffraction experiment provides insights into the dynamics of phase transitions and further allows to prepare and to characterize so far unknown transient phases of matter. This approach opens the door for further studies such as directly detecting spin, charge and orbital excitations in real time.²⁵ Improvements in beamline efficiency at FLASH and x-ray free-electron lasers that are presently coming online will widen the range of systems that can be studied considerably.

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